

Flatness Characterization: Two Approaches

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Flatness Characterization: Two Approaches

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Abstract. We survey two approaches to flatness necessary and sufficient conditions and compare them on examples.

1 Introduction

In this survey we consider underdetermined implicit systems of the form

$$F(x, \dot{x}) = 0 \tag{1}$$

with $x \in X$, X being an inifnitely differentiable manifold of dimension n, whose tangent bundle is denoted by TX, and $F: TX \to \mathbb{R}^{n-m}$ regular in the sense that $\operatorname{rk} \frac{\partial F}{\partial \dot{x}} = n - m$ in a suitable open dense subset of TX. Differential flatness, or more shortly, flatness was introduced in 1992 [20,11]. In the setting of implicit control systems it may be roughly described as follows: there exists a smooth mapping $x = \varphi(y, \dot{y}, \dots, y^{(r)})$ with $y = (y_1, \dots, y_m)^T$ of dimension m, $r = (r_1, \dots, r_m)^T \in \mathbb{N}^m$, such that

$$F(\varphi(y, \dot{y}, \dots, y^{(r)}), \dot{\varphi}(y, \dot{y}, \dots, y^{(r+1)})) \equiv 0$$
 (2)

with φ invertible in the sense that there exists a locally defined smooth mapping ψ and a multi-index s such that $y = \psi(x, \dot{x}, \dots, x^{(s)})$.

The vector y is called a *flat output*.

This concept has inspired an important literature. See [10,21,19,26,27,31] for surveys on flatness and its applications. Various formalisms have been introduced: finite dimensional differential geometric approaches [4,14,30], [32,28], differential algebra and related approaches [12,3,15], infinite dimensional differential geometry of jets and prolongations [13,33,19,6,7,23], [22,24], which is adopted here. The interested reader may refer to [1,13,16], [19,23,34] for more details.

The first part of the paper recalls the mathematical setting. In Section 3 the approch introduced in [19,2] for the characterization of differentially flat systems is recalled. Then, in Section 4, we introduce a novel characterization using the so-called Generalized Euler-Lagrange Operator. We conclude the paper with examples.

2 Implicit control systems on manifolds of jets of infinite order

Given an infinitely differentiable manifold X of dimension n, we denote its tangent space at $x \in X$ by $T_x X$, and its tangent bundle by TX. Let F be a meromorphic function from TX to \mathbb{R}^{n-m} . We consider an underdetermined implicit system of the form (1) regular in the sense that $\operatorname{rk} \frac{\partial F}{\partial \dot{x}} = n - m$ in a suitable open dense subset of TX.

Following [17,18], we consider the infinite dimensional manifold $\mathfrak X$ defined by $\mathfrak X \stackrel{\mathrm{def}}{=} X \times \mathbb R_\infty^n \stackrel{\mathrm{def}}{=} X \times \mathbb R^n \times \mathbb R^n \times \dots$, made of an infinite (but countable) number of copies of $\mathbb R^n$, with the global infinite set of coordinates $\overline{x} = \left(x, \dot{x}, \dots, x^{(k)}, \dots, \right)$, endowed with the product topology. Recall that, in this topology, a function φ from $\mathfrak X$ to $\mathbb R$ is continuous (resp. differentiable) if φ depends only on a finite (but otherwise arbitrary) number of variables and is continuous (resp. differentiable) with respect to these variables. C^∞ or analytic or meromorphic functions from $\mathfrak X$ to $\mathbb R$ are then defined as in the usual finite dimensional case since they only depend on a finite number of variables. We endow $\mathfrak X$ with the so-called trivial Cartan field ([16,34]) $\tau_{\mathfrak X} = \sum_{i=1}^n \sum_{j\geq 0} x_i^{(j+1)} \frac{\partial}{\partial x^{(j)}}$. We

so-called trivial Cartan field ([16,34]) $\tau_{\mathfrak{X}} = \sum_{i=1}^{n} \sum_{j\geq 0} x_{i}^{(j+1)} \frac{\partial}{\partial x_{i}^{(j)}}$. We also denote by $L_{\tau_{\mathfrak{X}}} \gamma = \sum_{i=1}^{n} \sum_{j\geq 0} x_{i}^{(j+1)} \frac{\partial \gamma}{\partial x_{i}^{(j)}} = \frac{d\gamma}{dt}$ the Lie derivative of a differentiable function γ along $\tau_{\mathfrak{X}}$ and $L_{\tau_{\mathfrak{X}}}^{k} \gamma$ its kth iterate. Since $\frac{d}{dt} x_{i}^{(j)} \stackrel{\text{def}}{=} \dot{x}_{i}^{(j)} = x_{i}^{(j+1)}$, the Cartan field acts on coordinates as a shift to the right. \mathfrak{X} is thus called manifold of jets of infinite order.

A regular implicit control system is defined as a triple $(\mathfrak{X}, \tau_{\mathfrak{X}}, F)$ with $\mathfrak{X} = X \times \mathbb{R}_{\infty}^n$, $\tau_{\mathfrak{X}}$ its associated trivial Cartan field, and F meromorphic from TX to \mathbb{R}^{n-m}) satisfying $\operatorname{rk} \frac{\partial F}{\partial \dot{x}} = n - m$ in a suitable open subset of TX.

We next consider the cotangent space $T_{\bar{x}}^*\mathfrak{X}$ with $dx_i^{(j)}$, $i=1,\ldots,n, j\geq 0$ as basis, dual to the $\frac{\partial}{\partial x_i^{(j)}}$'s. 1-forms on \mathfrak{X} are then defined in the usual way. The set of 1-forms is noted $\Lambda^1(\mathfrak{X})$. We also denote by $\Lambda^p(\mathfrak{X})$ the module of all the p-forms on \mathfrak{X} .

2.1 Flatness

We recall the following definitions and result [17,18,19]:

Given two regular implicit control systems $(\mathfrak{X}, \tau_{\mathfrak{X}}, F)$, with $\mathfrak{X} = X \times \mathbb{R}_{\infty}^{n}$, dim X = n and rk $\frac{\partial F}{\partial \dot{x}} = n - m$, and $(\mathfrak{Y}, \tau_{\mathfrak{Y}}, G)$, with $\mathfrak{Y} = Y \times \mathbb{R}_{\infty}^{p}$, dim Y = p, $\tau_{\mathfrak{Y}}$ its trivial Cartan field, and rk $\frac{\partial G}{\partial \dot{y}} = p - q$, we set $\mathfrak{X}_{0} = \{\overline{x} \in \mathfrak{X} | L_{\tau_{\mathfrak{X}}}^{k} F(\overline{x}) = 0, \ \forall k \geq 0\}$ and $\mathfrak{Y}_{0} = \{\overline{y} \in \mathfrak{Y} | L_{\tau_{\mathfrak{Y}}}^{k} G(\overline{y}) = 0, \ \forall k \geq 0\}$. They are endowed with the topologies and differentiable structures induced by \mathfrak{X} and \mathfrak{Y} respectively.

Definition 1 The control systems $(\mathfrak{X}, \tau_{\mathfrak{X}}, F)$ and $(\mathfrak{Y}, \tau_{\mathfrak{Y}}, G)$ are said locally Lie-Bäcklund equivalent (or shortly L-B equivalent) in a neighbourhood $\mathfrak{X}_0 \times \mathfrak{Y}_0$ of the pair $(\overline{x}_0, \overline{y}_0) \in \mathfrak{X}_0 \times \mathfrak{Y}_0$ if and only if

³ From now on, \overline{x} \overline{y} , ... stand for the sequences of jets of infinite order of x, y,...

- (i) there exists a one-to-one meromorphic mapping $\Phi = (\varphi, \dot{\varphi}, ...)$ from y_0 to x_0 satisfying $\Phi(\overline{y}_0) = \overline{x}_0$ and such that $\Phi_* \tau_{\mathfrak{Y}} = \tau_{\mathfrak{X}}$;
- (ii) there exists Ψ one-to-one and meromorphic from \mathfrak{X}_0 to \mathfrak{Y}_0 , with $\Psi = (\psi, \dot{\psi}, \ldots)$, such that $\Psi(\overline{x}_0) = \overline{y}_0$ and $\Psi_* \tau_{\mathfrak{X}} = \tau_{\mathfrak{Y}}$.

The mappings Φ and Ψ are called mutually inverse Lie-Bäcklund isomorphisms at $(\overline{x}_0, \overline{y}_0)$.

Definition 2 The implicit system $(\mathfrak{X}, \tau_{\mathfrak{X}}, F)$ is locally flat in a neighborhood of $(\overline{x}_0, \overline{y}_0) \in \mathfrak{X}_0 \times \mathbb{R}_{\infty}^m$ if and only if it is locally L-B equivalent around $(\overline{x}_0, \overline{y}_0)$ to the trivial implicit system $(\mathbb{R}_{\infty}^m, \tau_{\mathbb{R}_{\infty}^m}, 0)$. In this case, the mutually inverse L-B isomorphisms Φ and Ψ are called inverse trivializations.

Theorem 1 The system $(\mathfrak{X}, \tau_{\mathfrak{X}}, F)$ is locally flat at $(\overline{x}_0, \overline{y}_0) \in \mathfrak{X}_0 \times \mathbb{R}_\infty^m$ if and only if there exists a local meromorphic invertible mapping Φ from \mathbb{R}_∞^m to \mathfrak{X}_0 , with meromorphic inverse, satisfying $\Phi(\overline{y}_0) = \overline{x}_0$, and such that

$$\Phi^* dF = 0. (3)$$

3 Necessary and Sufficient Conditions: Generalized Moving Frame Structure Equations

3.1 Algebraic characterization of the differential of a trivialization

Consider the following matrix, polynomial with respect to the differential operator $\frac{d}{dt}$ (we use indifferently $\frac{d}{dt}$ for $L_{\tau_{\mathfrak{X}}}$ or $L_{\tau_{\mathbb{R}_{\infty}^{m}}}$, the context being unambiguous):

$$P(F) = \frac{\partial F}{\partial x} + \frac{\partial F}{\partial \dot{x}} \frac{d}{dt}, \quad P(\varphi) = \sum_{j \ge 0} \frac{\partial \varphi}{\partial y^{(j)}} \frac{d^j}{dt^j}$$
(4)

with P(F) (resp. $P(\varphi)$) of size $(n-m) \times n$ (resp. $n \times m$). Equation (3)) reads:

$$\Phi^* dF = P(F)P(\varphi)dy = 0. \tag{5}$$

Clearly, the entries of the matrices in (4) are polynomials in the differential operator $\frac{d}{dt}$ with meromorphic coefficients from $\mathfrak X$ to $\mathbb R$. We denote by $\mathfrak R$ the field of meromorphic functions from $\mathfrak X$ to $\mathbb R$ and by $\mathfrak R[\frac{d}{dt}]$ the (noncommutative) principal ideal ring of polynomials in $\frac{d}{dt}$ with coefficients in $\mathfrak R$. For $r,s\in\mathbb N$, let us denote by $\mathfrak M_{r,s}[\frac{d}{dt}]$ the module of $r\times s$ matrices over $\mathfrak R[\frac{d}{dt}]$ (see e.g. [8]). Matrices whose inverse belong to $\mathfrak M_{r,r}[\frac{d}{dt}]$ are called $unimodular\ matrices$. They form a multiplicative group denoted by $\mathfrak U_r[\frac{d}{dt}]$.

⁴ Note that if Φ is a meromorphic mapping from \mathfrak{Y} to \mathfrak{X} , the (backward) image by Φ of a 1-form is defined in the same way as in the finite dimensional context.

Every matrix $M \in \mathcal{M}_{r,s}[\frac{d}{dt}]$ admits a *Smith decomposition* (or diagonal reduction)

$$VMU = (\Delta, 0_{r,s-r}) \text{ if } r \leq s, \text{ and } \left(\frac{\Delta}{0_{r-s,s}}\right) \text{ if } s \leq r$$
 (6)

with $V \in \mathcal{U}_r[\frac{d}{dt}]$ and $U \in \mathcal{U}_s[\frac{d}{dt}]$ and Δ diagonal (see e.g. [8]). U and V are indeed non unique. We say that $U \in \mathsf{R} - \mathsf{Smith}(M)$ and $V \in \mathsf{L} - \mathsf{Smith}(M)$.

A matrix $M \in \mathcal{M}_{r,s}[\frac{d}{dt}]$ is said hyper-regular if and only if its Smith decomposition leads to $\Delta = I$. An interpretation of this property in terms of controllability in the sense of [9], may be found in [18].

From now on, we assume that P(F) is hyper-regular in a neighborhood of \overline{x}_0 . In place of (5), we first solve the matrix equation:

$$P(F)\Theta = 0 \tag{7}$$

where $\Theta \in \mathcal{M}_{n,m}[\frac{d}{dt}]$ is not supposed to be of the form $P(\varphi)$. It may be verified that matrices $\Theta \in \mathcal{M}_{n,m}[\frac{d}{dt}]$ satisfying (7) have the structure

$$\Theta = U \begin{pmatrix} 0_{n-m,m} \\ I_m \end{pmatrix} W \tag{8}$$

with $U \in \mathsf{R} - \mathsf{Smith}\,(P(F))$ and $W \in \mathcal{U}_m[\frac{d}{dt}]$ arbitrary. Clearly Θ is itself hyper-regular and admits the Smith decomposition

$$Q\Theta Z = QU \begin{pmatrix} 0_{n-m,m} \\ I_m \end{pmatrix} WZ = Q\hat{U}R = \begin{pmatrix} I_m \\ 0_{n-m,m} \end{pmatrix}$$
(9)

with
$$Q \in \mathcal{U}_n[\frac{d}{dt}], Z \in \mathcal{U}_m[\frac{d}{dt}], R = WZ$$
 and $\hat{U} = U\begin{pmatrix} 0_{n-m,m} \\ I_m \end{pmatrix}$.

3.2 Integrability

We denote by ω the m-dimensional vector 1-form defined by

$$\omega(\overline{x}) = \begin{pmatrix} \omega_1(\overline{x}) \\ \vdots \\ \omega_m(\overline{x}) \end{pmatrix} = (I_m, 0_{m,n-m}) Q(\overline{x}) dx_{|X_0}$$
 (10)

with Q given by (9), the restriction to \mathfrak{X}_0 meaning that $\overline{x} \in \mathfrak{X}_0$ satisfies $L^k_{\tau_{\mathfrak{X}}}F = 0$ for all k and that the $dx^{(k)}_j$ are such that $dL^k_{\tau_{\mathfrak{X}}}F = 0$ in \mathfrak{X}_0 for all k. Since Q is hyper-regular, the forms $\omega_1, \ldots, \omega_m$ are independent by construction.

Theorem 2 A necessary and sufficient condition for system (1) to be locally flat around $(\overline{x}_0, \overline{y}_0)$ is that there exist $U \in \mathsf{R} - \mathsf{Smith}(P(F))$, $Q \in \mathsf{L} - \mathsf{Smith}(\hat{U})$, with \hat{U} given by (9) and a matrix $M \in \mathfrak{U}_m[\frac{d}{dt}]$ such that $d(M\tau) = 0$.

We denote by $(\Lambda^p(\mathfrak{X}))^m$ the space of m-dimensional vector p-forms on \mathfrak{X} , by $(\Lambda(\mathfrak{X}))^m$ the space of m-dimensional vector forms of arbitrary degree on \mathfrak{X} , and by $\mathcal{L}_q((\Lambda(\mathfrak{X}))^m) = \bigcup_{p \geq 1} \mathcal{L}\left((\Lambda^p(\mathfrak{X}))^m, (\Lambda^{p+q}(\mathfrak{X}))^m\right)$ the space of linear operators from $(\Lambda^p(\mathfrak{X}))^m$ to $(\Lambda^{p+q}(\mathfrak{X}))^m$ for all $p \geq 1$, where $\mathcal{L}(\mathcal{P}, \mathcal{Q})$ denotes the set of linear mappings from a given space \mathcal{P} to a given space \mathcal{Q} .

In order to develop the expression $d(\mu\kappa)$ for $\mu \in \mathcal{L}_q((\Lambda(\mathfrak{X}))^m)$ and for all $\kappa \in (\Lambda^p(\mathfrak{X}))^m$ and all $p \geq 1$, we define the operator \mathfrak{d} by:

$$\mathfrak{d}(\mu) \kappa = d(\mu \kappa) - (-1)^q \mu \ d\kappa. \tag{11}$$

Note that (11) uniquely defines $\mathfrak{d}(\mu)$ as an element of $\mathcal{L}_{q+1}((\Lambda(\mathfrak{X}))^m)$.

Theorem 3 The system $(\mathfrak{X}, \tau_{\mathfrak{X}}, F)$ is locally flat iff there locally exists $\mu \in \mathcal{L}_1 ((\Lambda(\mathfrak{X}))^m)$, and a matrix $M \in \mathcal{U}_m[\frac{d}{dt}]$ such that

$$d\omega = \mu \ \omega, \qquad \mathfrak{d}(\mu) = \mu^2, \qquad \mathfrak{d}(M) = -M\mu.$$
 (12)

with the notation $\mu^2 = \mu\mu$ and where ω is defined by (10). In addition, if (12) holds true, a flat output y is obtained by integration of $dy = M\omega$.

Remark 1 Note that the two first conditions of (12) are comparable to conditions (A) and (B) of [6,7]. However, the last condition of (12) is different from condition (C) of [6,7] and is easier to check.

Note also that conditions (12) may be seen as a generalization in the framework of manifolds of jets of infinite order of Cartan's well-known moving frame structure equations (see e.g. [5]).

3.3 A Sequential Procedure

We start with P(F) hyper-regular and compute the vector 1-form ω defined by (10).

- 1. We identify the operator μ such that $d\omega = \mu\omega$ componentwise. It is proven in [19] that such μ always exists.
- 2. Among the possible μ 's, only those satisfying $\mathfrak{d}(\mu) = \mu^2$ are kept. It is shown in [19] that such μ always exists.
- 3. We then identify M such that $\mathfrak{d}(M) = -M\mu$ componentwise.
- 4. If, among such M's, there is a unimodular one, the system is flat and a flat output is obtained by integration of $dy = M\omega$. Otherwise the system is not flat.

More details and examples may be found in [18,19].

4 Necessary and Sufficient Conditions using the Generalized Euler-Lagrange Operator

Another way of analysing (3) consists in characterizing the change of coordinates corresponding to the mapping Φ in (3). More precisely (3) reads

$$\sum_{j=1}^{m} \sum_{k=0}^{r_j} \left(\frac{\partial F}{\partial x} \frac{\partial \varphi}{\partial y_j^{(k)}} dy_j^{(k)} + \frac{\partial F}{\partial \dot{x}} \frac{d}{dt} \left(\frac{\partial \varphi}{\partial y_j^{(k)}} \right) dy_j^{(k)} + \frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_j^{(k)}} dy_j^{(k+1)} \right) = 0$$
(13)

Since the one forms $dy_1, \ldots, dy_1^{(r_1)}, \ldots, dy_m, \ldots, dy_m^{(r_m)}$ are independent by assumption, (13) yields, for every $j = 1, \ldots, m$

$$\begin{cases} \frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(r_{j})}} = 0 \\ \frac{\partial F}{\partial x} \frac{\partial \varphi}{\partial y_{j}^{(k)}} + \frac{\partial F}{\partial \dot{x}} \frac{d}{dt} \left(\frac{\partial \varphi}{\partial y_{j}^{(k)}} \right) + \frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(k-1)}} = 0, \quad \forall k = 1, \dots, r_{j} \\ \frac{\partial F}{\partial x} \frac{\partial \varphi}{\partial y_{j}} + \frac{\partial F}{\partial \dot{x}} \frac{d}{dt} \left(\frac{\partial \varphi}{\partial y_{j}} \right) = 0 \end{cases}$$
(14)

The Generalized Euler-Lagrange operator \mathcal{E}_F associated to F is defined

$$\mathcal{E}_F = \frac{\partial F}{\partial x} - \frac{d}{dt} \left(\frac{\partial F}{\partial \dot{x}} \right) \tag{15}$$

In the case n-m=1, it is well-known that the curves that extremize the cost function $J=\int_0^T F(x,\dot{x})dt$ are those satisfying the Euler-Lagrange equation $\mathcal{E}_F = 0$, which justifies our terminology. Using (15) and elementary calculus, (14) yields:

Theorem 4 A necessary and sufficient condition for (1) to be diffferentially flat is that there exist (r_1, \ldots, r_m) with $\sum_{i=1}^m r_i + m \ge n$ and a solution φ of the following triangular system of PDEs in an open dense

$$\begin{cases}
\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(r_{j})}} = 0 \\
\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(l)}} = \sum_{k=0}^{r_{j}-l-1} (-1)^{k+1} \frac{d^{k}}{dt^{k}} \left(\mathcal{E}_{F} \frac{\partial \varphi}{\partial y_{j}^{(l+k+1)}} \right), \quad \forall l = 0, \dots, r_{j} - 1 , \\
0 = \sum_{k=0}^{r_{j}} (-1)^{k} \frac{d^{k}}{dt^{k}} \left(\mathcal{E}_{F} \frac{\partial \varphi}{\partial y_{j}^{(k)}} \right)
\end{cases} \tag{16}$$

satisfying $d\varphi_1 \wedge \ldots \wedge d\varphi_n \neq 0$.

Remark 2 If there exists a coordinate transformation φ that satisfies the conditions of Theorem 4 with given r_1, \ldots, r_m , meaning that the system is flat, then $g_j = \sum_{i=1}^n \frac{\partial \varphi_i}{\partial y_i^{(r_j)}} \frac{\partial}{\partial \dot{x}_i}$, if non zero, defines a ruled direction [32,25,19].

Examples 5

An Academic Example: Generalized Moving Frame Approach

We consider the 3-dimensional system with 2 inputs:

$$\dot{x}_1 = u_1, \ \dot{x}_2 = u_2, \ \dot{x}_3 = \sin\left(\frac{u_1}{u_2}\right)$$
 (17)

or, in implicit form:

$$F(x_1, x_2, x_3, \dot{x}_1, \dot{x}_2, \dot{x}_3) \triangleq \dot{x}_3 - \sin\left(\frac{\dot{x}_1}{\dot{x}_2}\right) = 0.$$
 (18)

It is readily seen that $P(F) = \left[-\cos(\frac{\dot{x}_1}{\dot{x}_2}) \dot{x}_2^{-1} \frac{d}{dt} \middle| \dot{x}_1 \cos(\frac{\dot{x}_1}{\dot{x}_2}) \dot{x}_2^{-2} \frac{d}{dt} \middle| \frac{d}{dt} \right]$ and that $VP(F)U = (1\ 0\ 0)$ with

$$V = 1, \quad U = \begin{pmatrix} \frac{\dot{x}_1}{a\dot{x}_2} & 1 + \frac{\dot{x}_1}{a(\dot{x}_2)^2} \cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right) \frac{d}{dt} & \frac{\dot{x}_1}{a\dot{x}_2} \frac{d}{dt} \\ \frac{1}{a} & \frac{1}{a\dot{x}_2} \cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right) \frac{d}{dt} & -\frac{1}{a} \frac{d}{dt} \\ 0 & 0 & 1 \end{pmatrix}$$
(19)

where $a = -\frac{1}{\dot{x}_2} \cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right) \left(\frac{\ddot{x}_1 \dot{x}_2 - \dot{x}_1 \ddot{x}_2}{(\dot{x}_2)^2}\right)$. Then, $Q\hat{U}R = \begin{pmatrix} 1 & 0 \\ 0 & 1 \\ 0 & 0 \end{pmatrix}$ is computed with

$$Q = \begin{pmatrix} 1 & -\frac{\dot{x}_1}{\dot{x}_2} & 0 \\ 0 & 1 \\ -\frac{1}{a\dot{x}_2}\cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right)\frac{d}{dt} & \frac{\dot{x}_1}{a(\dot{x}_2)^2}\cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right)\frac{d}{dt} & \frac{1}{a}\frac{d}{dt} \end{pmatrix}, \quad R = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
(20)

So,
$$(\omega_1 \ \omega_2)^T = \begin{pmatrix} 1 \ 0 \ 0 \end{pmatrix} Q dx = \begin{pmatrix} dx_1 - \frac{\dot{x}_1}{\dot{x}_2} dx_2 \ dx_3 \end{pmatrix}^T$$
 and $d\omega = \begin{pmatrix} \frac{1}{\sqrt{1-(\dot{x}_3)^2}} dx_2 \wedge dx_3 \ 0 \end{pmatrix}^T$. According to section 3.3, step 1,

$$\mu = \begin{pmatrix} 0 \left(-\frac{\dot{x}_3}{(1 - (\dot{x}_3)^2)^{\frac{3}{2}}} dx_2 \wedge d\dot{x}_3 + \eta d\dot{x}_3 \right) \wedge \frac{d}{dt} \\ 0 & 0 \end{pmatrix}. \tag{21}$$

Step 2 yields $\eta = \frac{x_2 \dot{x}_3}{(1-\dot{x}_3)^{\frac{3}{2}}} + \sigma(\dot{x}_3)$. For step 3 we set $M = \begin{pmatrix} 1 & m_{12} \frac{d}{dt} \\ 0 & 1 \end{pmatrix}$ which yields $m_{12} = -\left(\frac{x_2}{\sqrt{1-(\dot{x}_3)^2}} + \sigma_1(\dot{x}_3)\right)$ with σ_1 a primitive of σ . Thus, $d(M\omega) = 0$ and setting $(dy_1 \ dy_2)^T = M\omega$, one obtains

$$y_1 = x_1 - \frac{\dot{x}_1}{\dot{x}_2} x_2 + \sigma_2(\dot{x}_3), \quad y_2 = x_3$$
 (22)

where $\sigma_2(\dot{x}_3)$ is an arbitrary meromorphic function (a primitive of σ_1). By inversion of (22) we get

$$x_{1} = y_{1} - \arcsin(\dot{y}_{2}) \frac{\sqrt{1 - (\dot{y}_{2})^{2}}}{\ddot{y}_{2}} (\dot{y}_{1} - \sigma_{1}(\dot{y}_{2})\ddot{y}_{2}) - \sigma_{2}(\dot{y}_{2})$$

$$x_{2} = -\frac{\sqrt{1 - (\dot{y}_{2})^{2}}}{\ddot{y}_{2}} (\dot{y}_{1} - \sigma_{1}(\dot{y}_{2})\ddot{y}_{2})$$

$$x_{3} = y_{2}$$

$$(23)$$

Academic Example: Euler-Lagrange Operator

We consider once more the example (18). We have

$$\frac{\partial F}{\partial \dot{x}} = \left(-\dot{x}_2^{-1}\cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right), \dot{x}_1\dot{x}_2^{-2}\cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right), 1\right), \ \mathcal{E}_F = (\eta_1, \eta_2, 0) \quad (24)$$

with
$$\eta_1 = -\frac{\ddot{x}_2}{\dot{x}_2^2} \cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right) - \frac{\ddot{x}_1\dot{x}_2 - \dot{x}_1\ddot{x}_2}{\dot{x}_2^3} \sin\left(\frac{\dot{x}_1}{\dot{x}_2}\right)$$
 and $\eta_2 = -\frac{\ddot{x}_1\dot{x}_2 - 2\dot{x}_1\ddot{x}_2}{\dot{x}_2^3} \cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right) + \frac{\dot{x}_1(\ddot{x}_1\dot{x}_2 - \dot{x}_1\ddot{x}_2)}{\dot{x}_2^4} \sin\left(\frac{\dot{x}_1}{\dot{x}_2}\right)$. The first two equations of (16), with $r_1 = r_2 = 2$, read

$$-\frac{1}{\dot{x}_2}\cos\left(\frac{\dot{x}_1}{\dot{x}_2}\right)\left(\frac{\partial\varphi_1}{\partial\dot{y}_j} - \frac{\dot{x}_1}{\dot{x}_2}\frac{\partial\varphi_2}{\partial\dot{y}_j}\right) + \frac{\partial\varphi_3}{\partial\dot{y}_j} = 0, \quad j = 1, 2$$
 (25)

If we assume that $\frac{\partial \varphi_3}{\partial \ddot{y}_i} = \frac{\partial \varphi_3}{\partial \ddot{y}_i} = 0$, j=1,2 and introduce the variable

$$\psi = \frac{\dot{x}_1}{\dot{x}_2} \tag{26}$$

with $\frac{\partial}{\partial \dot{u}} \psi = 0$ we obtain from (25)

$$\frac{\partial \varphi_1}{\partial \ddot{y}_j} - \psi \frac{\partial \varphi_2}{\partial \ddot{y}_j} = \frac{\partial}{\partial \ddot{y}_j} (\varphi_1 - \psi \varphi_2) = 0, \quad j = 1, 2$$

Setting $\kappa(y, \dot{y}) = \varphi_1 - \psi \varphi_2$, we get

$$\dot{\kappa} = \dot{\varphi}_1 - \psi \dot{\varphi}_2 - \dot{\psi} \varphi_2 = -\dot{\psi} \varphi_2 \tag{27}$$

Using the definition of κ and (27) we obtain:

$$\varphi_1 = \kappa - \frac{\dot{\kappa}\sqrt{1 - \dot{\varphi}_3}}{\ddot{\varphi}_3} \arcsin(\dot{\varphi}_3), \quad \varphi_2 = -\frac{\dot{\kappa}}{\ddot{\varphi}_3} \sqrt{1 - \dot{\varphi}_3}, \quad \varphi_3 = \varphi_3(y)$$
(28)

Choosing $\varphi_3 = y_2$, $\kappa = y_1$, we arrive at the invertible transformation

$$x_1 = \varphi_1 = y_1 - \frac{\dot{y}_1}{\ddot{y}_2} \sqrt{1 - \dot{y}_2^2} \arcsin(\dot{y}_2), \ x_2 = \varphi_2 = -\frac{\dot{y}_1}{\ddot{y}_2} \sqrt{1 - \dot{y}_2^2},$$

with $x_3 = \varphi_3 = y_2$, which gives the same formula as (23) with $\sigma_1 =$ $\sigma_2 = 0$. Hence (y_1, y_2) is indeed a flat output, which implies that the remaining equations of (16) are satisfied.

An Example Proposed by P. Rouchon

Consider the implicit control system

$$F(x, \dot{x}) = \dot{x}_1 \dot{x}_3 - (\dot{x}_2)^2 = 0.$$
 (29)

We thus have $\frac{\partial F}{\partial x} = \begin{pmatrix} 0 & 0 & 0 \end{pmatrix}$, $\frac{\partial F}{\partial \dot{x}} = \begin{pmatrix} \dot{x}_3 & -2\dot{x}_2 & \dot{x}_1 \end{pmatrix}$ and

$$\mathcal{E}_F = \frac{\partial F}{\partial x} - \frac{d}{dt} \left(\frac{\partial F}{\partial \dot{x}} \right) = -\frac{d}{dt} \left(\frac{\partial F}{\partial \dot{x}} \right) = \left(-\ddot{x}_3 \ 2\ddot{x}_2 \ -\ddot{x}_1 \right).$$

The lowest possible choice of (r_1, r_2) in Theorem 4 is $r_1 = r_2 = 1$. However, there is no solution of (16) for these values, and we choose $r_1 = r_2 = 2$. The two first equations of (16) read

$$\dot{\varphi}_3 \frac{\partial \varphi_1}{\partial \ddot{y}_i} - 2\dot{\varphi}_2 \frac{\partial \varphi_2}{\partial \ddot{y}_i} + \dot{\varphi}_1 \frac{\partial \varphi_3}{\partial \ddot{y}_i} = 0, \quad j = 1, 2$$
 (30)

We divide (30) by $\dot{\varphi}_3$ to obtain

$$\frac{\partial \varphi_1}{\partial \ddot{y}_j} - 2\psi \frac{\partial \varphi_2}{\partial \ddot{y}_j} + \psi^2 \frac{\partial \varphi_3}{\partial \ddot{y}_j} = 0, \quad j = 1, 2$$
(31)

where, taking account of the system equation (29),

$$\psi = \frac{\dot{\varphi}_2}{\dot{\varphi}_3} = \sqrt{\frac{\dot{\varphi}_1}{\dot{\varphi}_3}}.$$
 (32)

If we assume that ψ doesn't depend on \ddot{y}_1 and \ddot{y}_2 , equation (31) reads $\frac{\partial}{\partial \ddot{y}_j} \left(\varphi_1 - 2\psi \varphi_2 + \psi^2 \varphi_3 \right) = 0$, for j = 1, 2. In other words, there exists a function κ satisfying $\frac{\partial \kappa}{\partial \ddot{y}_j} = 0$ for j = 1, 2, such that

$$\varphi_1 - 2\psi\varphi_2 + \psi^2\varphi_3 = \kappa \tag{33}$$

Differentiating the latter relation with respect to t, and taking into account the relation $\dot{\varphi}_1 - 2\psi\dot{\varphi}_2 + \psi^2\dot{\varphi}_3 = 0$ obtained from (29) and (32), we get

$$\varphi_2 - \psi \varphi_3 = -\frac{\dot{\kappa}}{2\dot{\psi}}.\tag{34}$$

We again differentiate the latter relation with respect to t to obtain

$$\varphi_3 = \frac{\ddot{\kappa}\dot{\psi} - \dot{\kappa}\ddot{\psi}}{2\dot{\psi}^3} \tag{35}$$

thanks to $\dot{\varphi}_2 - \psi \dot{\varphi}_3 = 0$ from (32). Thus, solving the system (33)–(35), we immediately obtain

$$\varphi_{1} = \kappa - \psi \frac{\dot{\kappa}}{\dot{\psi}} + \psi^{2} \left(\frac{\ddot{\kappa} \dot{\psi} - \dot{\kappa} \ddot{\psi}}{2 \dot{\psi}^{3}} \right)
\varphi_{2} = -\frac{\dot{\kappa}}{2 \dot{\psi}} + \psi \left(\frac{\ddot{\kappa} \dot{\psi} - \dot{\kappa} \ddot{\psi}}{2 \dot{\psi}^{3}} \right)
\varphi_{3} = \frac{\ddot{\kappa} \dot{\psi} - \dot{\kappa} \ddot{\psi}}{2 \dot{\psi}^{3}}$$
(36)

where κ and ψ are arbitrary functions of $y_1, y_2, \dot{y}_1, \dot{y}_2$.

Note that choosing $\kappa = y_1$ and $\psi = y_2$ yields, after inversion of (36) with (32):

$$y_1 = x_1 - 2x_2 \frac{\dot{x}_2}{\dot{x}_3} + x_3 \frac{\dot{x}_1}{\dot{x}_3}, \quad y_2 = \frac{\dot{x}_2}{\dot{x}_3},$$

which is similar to the solution obtained by F. Ollivier⁵.

Similarly, the solution of K. Schlacher and M. Schöberl [29] may be

⁵ personal communication

recovered by posing $\kappa = y_1 - y_2 \frac{\dot{y}_1}{\dot{y}_2}$ and $\psi = \frac{\dot{y}_1}{2\dot{y}_2}$ which, again after inversion of (36) with (32), yields:

$$y_1 = x_1 - x_3 \frac{\dot{x}_1}{\dot{x}_3}, \quad y_2 = x_2 - x_3 \frac{\dot{x}_2}{\dot{x}_3}.$$

6 Conclusion

In this survey we presented two dual approaches to flatness necessary and sufficient conditions, one based on the integration of 1-forms and the second based on the integration of a set of PDEs involving a generalized Euler-Lagrange operator. Their complexity is compared on examples.

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